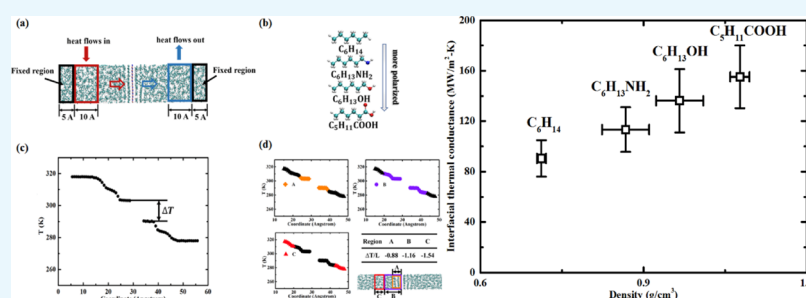


# 1 Role of Molecular Polarity in Thermal Transport of Boron Nitride– 2 Organic Molecule Composites

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7 **ABSTRACT:** Understanding the role of fillers in the thermal transport of composite materials is of great importance to  
8 engineering better materials. The filler induces material interfaces within the composite, which influence the thermal transport  
9 between the matrix and themselves. The filler can also alter the molecular arrangement of the matrix in its vicinity, which may  
10 also impact the thermal transport ability. In this paper, molecular dynamics simulations are performed to study the thermal  
11 transport across the matrix–filler interfaces in hexagonal boron nitride (h-BN)–organic molecule composites. Four different  
12 organic molecules are studied as the matrixes. They include hexane (C<sub>6</sub>H<sub>14</sub>), hexanamine (C<sub>6</sub>H<sub>13</sub>NH<sub>2</sub>), hexanol (C<sub>6</sub>H<sub>13</sub>OH),  
13 and hexanoic acid (C<sub>5</sub>H<sub>11</sub>COOH), which feature the same molecular backbone but increasingly different polar functional  
14 groups. The nominal local thermal conductivities of the hexane matrix with varying distances to the interface are calculated to  
15 demonstrate the influence of the filler on the thermal transport properties of the matrix. It is found that a more polar matrix  
16 exhibits a higher density in the near-interface region and a higher nominal local thermal conductivity, suggesting that the  
17 interfacial interaction can impact the local heat transfer ability of the matrix. In addition, the more polar matrix also leads to a  
18 larger interfacial thermal conductance with h-BN (hexane: 90.47 ± 14.49 MW/m<sup>2</sup> K, hexanamine: 113.38 ± 17.72 MW/m<sup>2</sup> K,  
19 hexanol: 136.16 ± 25.12 MW/m<sup>2</sup> K, and hexanoic acid: 155.17 ± 24.89 MW/m<sup>2</sup> K) because of the higher matrix density near  
20 the interface and thus more atoms exchanging energy with the filler. The results of this study may provide useful information for  
21 designing composite materials for heat transfer applications.

## 22 ■ INTRODUCTION

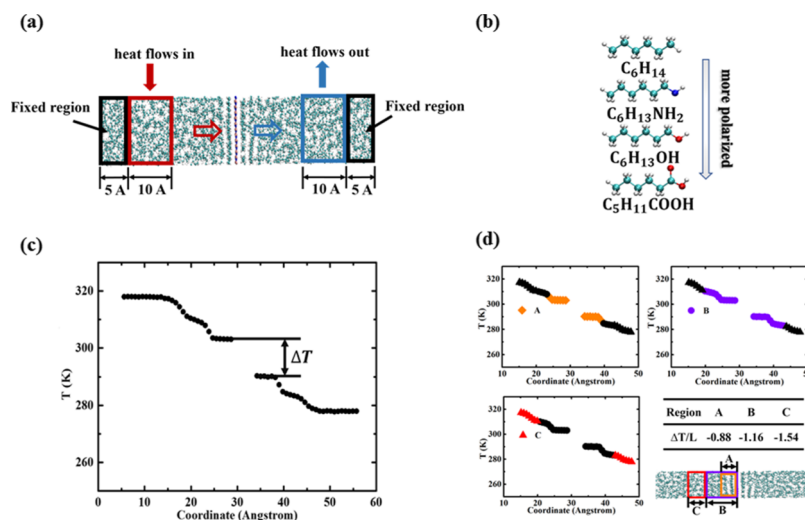
23 Thermal transport is critical to a wide range of applications  
24 such as advanced electronics,<sup>1,2</sup> optoelectronics, photovoltaic  
25 solar cells,<sup>3</sup> and Li-ion batteries.<sup>4</sup> It is also the key to  
26 determining the lifetime and performance of these devices.  
27 There has been an increasing demand for polymeric materials  
28 with high thermal conductivity that can dissipate heat  
29 generated by electronic devices in operation. To create  
30 efficient heat transfer pathways, various attempts have been  
31 exploited to increase the thermal conductivity of these  
32 materials by compositing high thermal conductivity fillers,  
33 such as carbonaceous materials, boron nitride, and silicon  
34 nitride.<sup>5–15</sup> Thus, to facilitate the design of nanocomposites  
35 with high thermal conductivity, a fundamental and compre-  
36 hensive understanding of the thermal transport in polymeric  
37 nanocomposites is essential.

38 Though it is universally known that adding high thermal  
39 conductivity materials into a low thermal conductivity matrix  
40 can improve the thermal transport,<sup>16,17</sup> the fundamental

41 mechanism that links the role of fillers to the overall thermal  
42 transport in nanocomposites is still lacking. A typical problem  
43 being addressed in nanocomposites is the thermal transport  
44 across the interface between the matrix and the fillers.<sup>18,19</sup>  
45 Extensive studies have demonstrated that the interface plays a  
46 critical role in determining the thermal energy transport in  
47 nanocomposites, as the interface can influence and even  
48 dominate the thermal transport in nanocomposites.<sup>16,17</sup> Recent  
49 studies show that the interfacial thermal resistance could be  
50 decreased by tuning the alignment or surface functionalization  
51 of graphene in composites.<sup>18,20</sup> Most of these studies only  
52 address the heat flux exchange across the interface. Beyond the  
53 direct impact on interfacial thermal conductance, the fillers  
54 also affect the nominal local thermal conductivity of the matrix,  
55 which has been studied much less intensively. It has been

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**Figure 1.** (a) Simulation setup for ITC calculations in the NEMD method: heat flows across the interface from the heat source (red region) to the heat sink (blue region), with the fixed regions (black region) set at the end of the system; the arrows in the graph represent the direction of heat flow; the length of both the heat source and heat sink is 10 Å, and the length of each fixed region is 5 Å. (b) The chemical structures of hexane ( $C_6H_{14}$ ), hexanamine ( $C_6H_{13}NH_2$ ), hexanol ( $C_6H_{13}OH$ ), and hexanoic acid ( $C_5H_{11}COOH$ ); the order of polarity is  $C_5H_{11}COOH > C_6H_{13}OH > C_6H_{13}NH_2 > C_6H_{14}$ . (c) Steady-state temperature profile of the simulation system (shown in a) and the temperature difference ( $\Delta T$ ) across the interface. (d) Temperature points leveraged to calculate  $\Delta T$  for regions A, B, and C (highlighted in different colors), where  $\Delta T$  is the temperature difference between the two ends of the chosen region and  $\Delta T/L$  is the temperature difference divided by the region length in each region; the smaller the absolute value of  $\Delta T/L$ , the larger the nominal local thermal conductivity.

demonstrated that the arrangement of matrix molecules close to the interface is influenced by their dipoles,<sup>21,22</sup> resulting in a different density near the interface.<sup>23,24</sup> This may imply that the nominal local thermal properties of the matrix can also be changed by the fillers, as the thermal conductivity of polymers is a strong function of their molecular-level conformations.<sup>25,26</sup> In this work, we perform molecular dynamics (MD) simulations and analyses to explore the impact of the h-BN fillers on the thermal transport. As an insulator, h-BN fillers are preferred over graphene, a semimetal, in many applications (e.g., electronics packaging). The size dependence of interfacial thermal conductance is first examined on a single-layer hexagonal boron nitride (h-BN)/hexane ( $C_6H_{14}$ ) system. Then, a series of h-BN interfaces with organic molecules of different polarization groups as the matrixes are investigated. These include hexanamine ( $C_6H_{13}NH_2$ ), hexanol ( $C_6H_{13}OH$ ), and hexanoic acid ( $C_5H_{11}COOH$ ). We note that our focus of this study is to understand how different functional groups in organic molecules influence thermal transport across their interfaces with h-BN. The selection of the above four materials largely eliminates the impact of the molecular chain length and backbone on the interfacial thermal transport, leaving the functional group the only factor that contributes to the change in interfacial thermal conductance. First, we define matrix regions based on the distances to the interface and calculate the nominal local thermal conductivities of these regions. It is interesting to find that the near-interface region tends to exhibit a higher density, which leads to a higher nominal local thermal conductivity. Compared to nonpolar matrix molecules, polar molecules are more strongly attracted to the BN layer and form a near-interface region of both higher density and higher nominal local thermal conductivity. Because of the higher near-interface density, more atoms join the energy exchange across the interface and lead to a higher interfacial thermal conductance. These results, although fundamental in

nature, may have implications in the design of materials or nanostructures for different applications.

## SIMULATION METHODS

Here, single-layer h-BN is used as the filler material, and four types of organic molecules, hexane ( $C_6H_{14}$ ), hexanamine ( $C_6H_{13}NH_2$ ), hexanol ( $C_6H_{13}OH$ ), and hexanoic acid ( $C_5H_{11}COOH$ ) are studied as matrixes. With similar carbon backbones and different end groups, the four types of organic molecules should have similar vibrational spectra,<sup>27</sup> but different polarizations. According to the group electronegativity of the functional groups in the four matrices ( $CH_3$  2.17;  $NH_2$  2.39;  $OH$  2.85;  $COOH$  3.09), the order of their polarity is  $C_5H_{11}COOH > C_6H_{13}OH > C_6H_{13}NH_2 > C_6H_{14}$  (Figure 1b).<sup>28</sup> Thus, the differences in the interfacial thermal conductance of these four h-BN–organic molecule systems may be mostly attributed to polarization.

The Tersoff potential<sup>29</sup> is used to describe the interaction among the h-BN atoms, and a modified universal force field<sup>30</sup> is used for the nonbonding interaction between h-BN and organic molecules. All 12-6 Lennard-Jones (L-J) coefficients associated with the h-BN atoms are listed in Table 1, and the cutoff is set to 10 Å. The charges of the B and N atoms in h-

**Table 1.** L-J Potential Parameters between the h-BN Atoms and the Rest of the Atoms

pair type	$\epsilon$ (kcal/mol)	$\sigma$ (Å)
C–B	0.1374772708486750	3.53419521486273
O–B	0.1039230480000000	3.37784249000000
N–B	0.1114450530000000	3.44911438700000
H–B	0.0889943818451480	3.10433658336001
C–N	0.0851175657546667	3.34577013597604
O–B	0.0643428310000000	3.18941741100000
N–N	0.0690000000000000	3.66000000000000
H–N	0.0550999092558236	2.91591150447333

113 BN are respectively 1.0 and  $-1.0$ , and the Hockney particle–  
 114 particle particle–mesh<sup>31</sup> method is used here for the  
 115 evaluation of coulomb energies and forces. The organic  
 116 molecules are simulated using the polymer consistent force  
 117 field.<sup>32</sup> All simulations are carried out using the large-scale  
 118 atomic/molecular massively parallel simulator.<sup>33</sup> A timestep  
 119 size of 0.25 fs is used for all simulations.

120 A typical structure is shown in Figure 1a. A single layer of h-  
 121 BN is centered in the simulation domain with periodic  
 122 boundary conditions (PBCs) in all three directions. The PBCs  
 123 in the lateral directions model h-BN effectively infinite in size  
 124 without edges. This is reasonable as the focus here is the  
 125 interfaces between the basal plane of h-BN and the organic  
 126 matrixes, and in reality, such interfaces are dominant in the  
 127 composite. Each organic molecular system contains 200  
 128 molecules. The whole system is heated up to 600 K under  
 129 an *NPT* (constant number of atoms, pressure, and temper-  
 130 ature) ensemble to achieve a disordered amorphous phase and  
 131 then cooled down to 300 K with an annealing speed of 12 K/  
 132 ps. We note that the PBC in the direction perpendicular to the  
 133 h-BN surface is necessary to ensure that the *NPT* optimization  
 134 leads to the correct density of the organic liquid, which is  
 135 important to the interfacial thermal transport.<sup>34,35</sup>

136 Nonequilibrium MD (NEMD) is used to calculate the  
 137 interfacial thermal conductance. After the *NPT* relaxation at  
 138 300 K and 1 atm for 1.5–2 ns, the system is then simulated  
 139 using an *NVE* (constant number of atoms, volume, and  
 140 energy) ensemble with a heat source (320 K) and sink (280 K)  
 141 imposed at the two ends of the system using Langevin  
 142 thermostats (Figure 1a). A layer of atoms at each end of the  
 143 simulation domain is fixed to prevent the heat transfer across  
 144 the periodic boundaries, forcing all heat flux to cross the h-  
 145 BN–organic molecule interface. The fixed atoms also prevent  
 146 the translational drift of the whole system and thus help “lock”  
 147 the position of the interface and extract the temperature  
 148 profile. The relatively large temperature difference (40 K) is  
 149 used to establish a measurable temperature gap across the  
 150 interface. Except the fixed atoms, the rest of the system,  
 151 including the thermostated regions, is simulated in the *NVE*  
 152 ensemble for 10 ns for nominal local thermal conductivity and  
 153 interfacial thermal conductance calculations at steady state. A  
 154 typical temperature profile at the steady state is shown in  
 155 Figure 1c. The heat flux ( $q$ ) is calculated by averaging the  
 156 energy input and output rates from the heat source and sink.  
 157 The temperature jump ( $\Delta T$ ) across the interface is defined by  
 158 the temperature difference between the organic molecules at  
 159 the two sides of the h-BN layer. The interfacial thermal  
 160 conductance ( $G$ ) is then calculated as  $G = q/\Delta T$ . The nominal  
 161 local thermal conductivity ( $\kappa$ ) of the matrix is calculated as  $\kappa =$   
 162  $q/(\Delta T/L)$ , where  $\Delta T$  is the temperature difference between  
 163 the two ends of the chosen region and  $L$  is the length of the  
 164 chosen region. Figure 1d shows the temperature points used to  
 165 calculate  $\Delta T$  in regions A, B, and C; h-BN/hexane system is  
 166 used here as an example. The whole system is divided into  
 167 several equal-length bins, and the corresponding temperature is  
 168 calculated for each bin. Then, the temperature points selected  
 169 for each region are strictly based on the length of each region.  
 170 The nominal local thermal conductivity is used to characterize  
 171 the heat transfer capability of the region near the interface,  
 172 which cannot be interpreted as the thermal conductivity of the  
 173 organic liquid. We conduct two independent NEMD  
 174 simulations for each system, and in each simulation, data at  
 175 eight different time intervals (1 ns each) in the steady state are

used for property calculations. Therefore, there are 16 data  
 176 points averaged for each value of the interfacial thermal  
 177 conductance. The error bars are the standard deviations of  
 178 these data.  
 179

## RESULTS AND DISCUSSION

The test on the size dependence of interfacial thermal  
 181 conductance is first conducted to choose the appropriate  
 182 system size for later simulations. The interfacial thermal  
 183 conductance of systems with 200, 400, and 600 hexane  
 184 molecules are respectively  $88.01 \pm 13.77$ ,  $104.21 \pm 14.49$ , and  
 185  $93.62 \pm 32.43$  MW/m<sup>2</sup> K. As all results are within their error  
 186 bars, no size dependence is observed, and thus the system with  
 187 the size of 200 organic molecules is selected, so as to save the  
 188 simulation time but still produce justified calculations. We note  
 189 that the thermal conductance for the h-BN/hexane interface  
 190 calculated here is larger compared to that of the graphene/  
 191 polyethylene interfaces ( $61$  MW/m<sup>2</sup> K),<sup>34</sup> and this may be  
 192 attributed to the stronger L-J interactions and additional  
 193 electrostatic interactions across the h-BN/hexane interface.  
 194

Next, the nominal local thermal conductivities and densities  
 195 of the hexane matrix with varying distances to the interface are  
 196 calculated, and the results are shown in Figure 2. Three regions  
 197 are

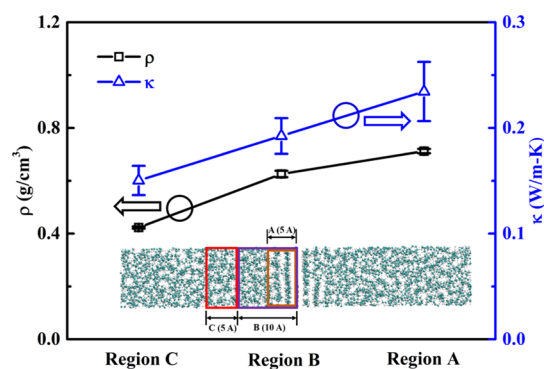
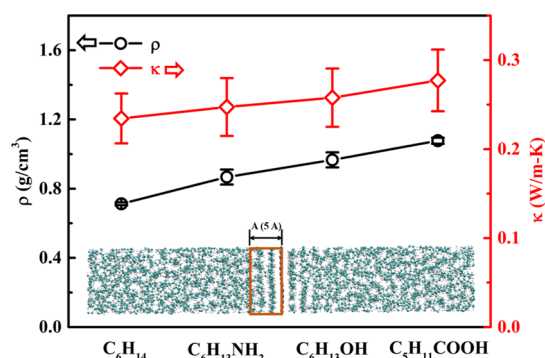


Figure 2. Nominal local thermal conductivities and densities of different regions in the h-BN/hexane system.

in the hexane matrix are chosen for exploration. Region A is  
 198 the region within a distance of 5 Å next to the interface. Region  
 199 B is the region within a distance of 10 Å next to the interface.  
 200 Region C is the whole system, except for region B, region A,  
 201 fixed region, heat source, and heat sink. Further, all the three  
 202 regions contain only organic molecules. In Figure 2, it is found  
 203 that when the hexane matrix gets closer to the h-BN interface,  
 204 its nominal local thermal conductivity becomes larger. The  
 205 higher nominal local thermal conductivity of the matrix near  
 206 the interface can be attributed to the higher density of organic  
 207 molecules,<sup>36</sup> as the corresponding density increases when  
 208 approaching the interface.  
 209

Figure 3 shows the thermal conductivities and densities of  
 210 different polar matrixes in region A, which is also called the  
 211 near-interface region. According to the results in Figure 2, the  
 212 region A shows the highest density and largest nominal local  
 213 thermal conductivity. Therefore, the region A is being chosen  
 214 for comparison among differently polarized matrixes, so as to  
 215 offer the most representative results. It is interesting to find  
 216 that when the organic matrix becomes more polarized, the  
 217 organic molecules in the near-interface region have larger  
 218 nominal local thermal conductivity, which could be explained  
 219 by the corresponding higher density, as it is known that a  
 220

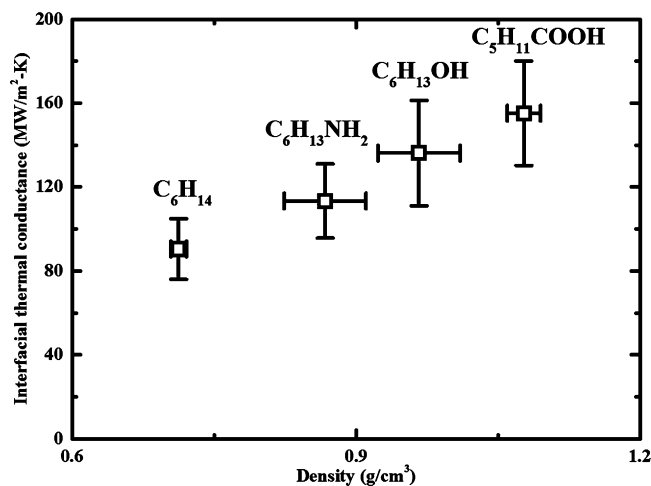




**Figure 3.** Nominal local thermal conductivities and densities of organic molecules at the near-interface region in the h-BN/hexane, h-BN/hexanamine, h-BN/hexanol, and h-BN/hexanoic acid systems.

221 higher density of a liquid can lead to a higher nominal local  
222 thermal conductivity because of the closer intermolecular  
223 distance.<sup>26</sup>

224 Finally, the thermal transport across the interface is being  
225 explored via calculating the interfacial thermal conductance of  
226 all the matrixes mentioned above. As Figure 4 shows, when the



**Figure 4.** Interfacial thermal conductance of h-BN/hexane, h-BN/hexanamine, h-BN/hexanol, and h-BN/hexanoic acid as a function of the density of the organic molecules at the near-interface region.

227 matrix becomes more polarized, the interfacial thermal  
228 conductance becomes larger. Such an observation is well-  
229 correlated to the increasing trend of the local density of the  
230 matrix near the interface as the molecules become more polar.  
231 The more polar organic molecules are attracted by h-BN closer  
232 to the interface because of the stronger electrostatic  
233 interactions. Such a shorter interatomic distance can  
234 significantly enhance the heat flux contributed by the L-J  
235 interaction. In addition, the higher local density near the  
236 interface also allows more organic molecular atoms to have  
237 such close “contact” with the h-BN layer. In other words, it is  
238 the collaborative effect of the electrostatic and L-J interactions  
239 that leads to the observed increase in interfacial thermal  
240 conductance as a function of matrix molecular polarity.

## 241 ■ CONCLUSIONS

242 In summary, we use the steady-state NEMD method to explore  
243 both the thermal transport across the interface and heat

transfer within the matrix. Four types of organic molecules 244  
with different polarities are studied as the matrixes (hexane 245  
( $C_6H_{14}$ ), hexanamine ( $C_6H_{13}NH_2$ ), hexanol ( $C_6H_{13}OH$ ), and 246  
hexanoic acid ( $C_5H_{11}COOH$ ). The nominal local thermal 247  
conductivities of the hexane matrix with varying distances to 248  
the interface are calculated to demonstrate the influence of the 249  
fillers on the matrix. It is found that the near-interface region 250  
will form a denser layer and exhibit a higher nominal local 251  
thermal conductivity, and such an effort is more pronounced 252  
for a more polar matrix. The higher density of the near- 253  
interface region in the more polar matrix also leads to a larger 254  
interfacial thermal conductance (hexane:  $90.47 \pm 14.49$  MW/ 255  
 $m^2$  K, hexanamine:  $113.38 \pm 17.72$  MW/ $m^2$  K, hexanol: 256  
 $136.16 \pm 25.12$  MW/ $m^2$  K, and hexanoic acid:  $155.17 \pm 24.89$  257  
MW/ $m^2$  K). These conclusions may be generalized to other 258  
organic molecules and polar filler materials. The results of this 259  
study may provide useful information for designing composite 260  
materials for heat transfer applications. 261

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### 272 Notes

273 The authors declare no competing financial interest.

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